

シミュレーション工学 (後半)

第3回 線形静弾性解析

線形と非線形

- 現実是非線形
- 実用上は線形でOKなことがほとんど
 - 微小変形, 固有振動
 - CAEとして製品開発に使われている解析の90%は線形解析
- 非線形でなければ表現できない現象
 - 座屈, 破壊, 接触

線形解析とは

- 線形の偏微分方程式
↓
- 線形の連立方程式

- 荷重が2倍になれば変位, 応力も2倍
- 複数の荷重は重ね合わせればよい

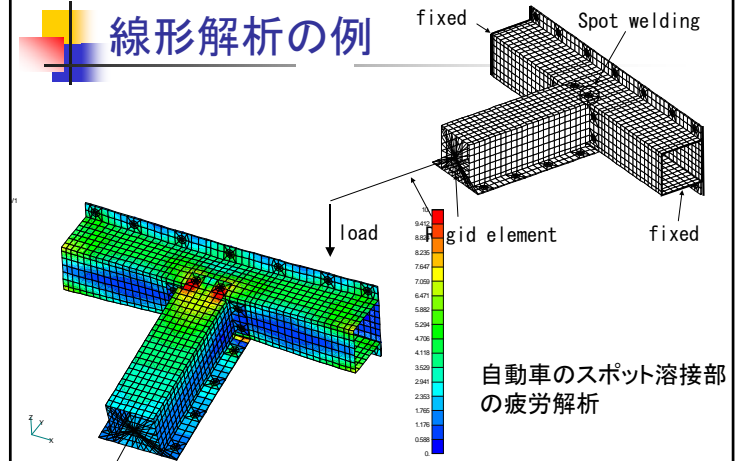
線形静弾性解析で扱える範囲

- 静的問題
 - 時間変化なし
- 微小変形
 - 変形が小さい
- 弾性変形
 - 塑性する前まで

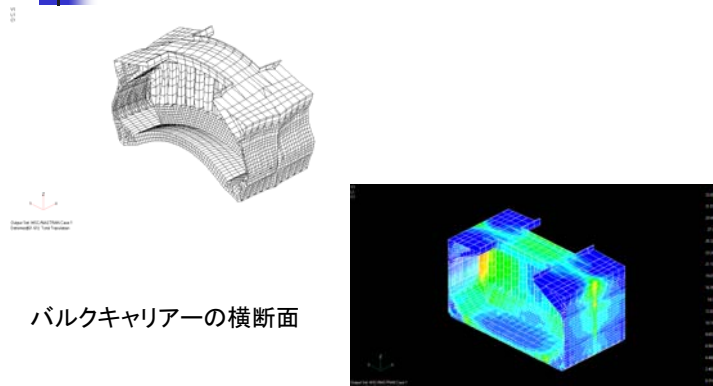
なぜ実用上線形解析が多いか

- 非線形解析は計算時間がかかる
 - 線形の数倍
- 非線形解析には専門知識, ノウハウが不可欠
- 応力を評価すれば破壊が予測可能
 - 実際のメカニズムはもっと複雑. 亀裂進展など
- 構造物は普通は固いので変形は微少

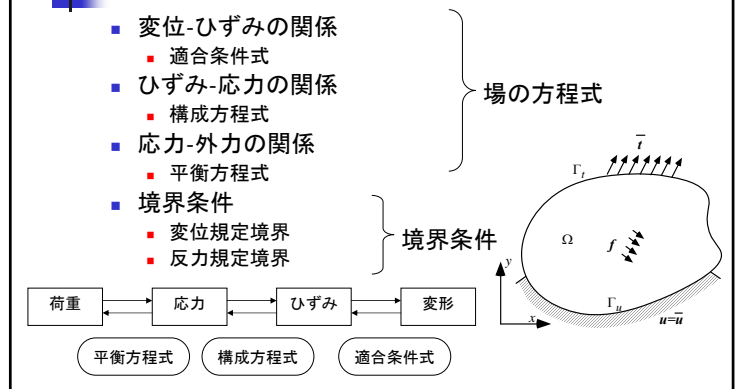
線形解析の例



線形解析の例2



固体力学の基礎方程式

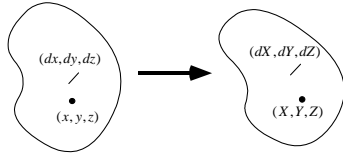


変位

変形前、変形後の座標

$$\begin{cases} u = X - x \\ v = Y - y \\ w = Z - z \end{cases}$$

$$ds^2 = dx^2 + dy^2 + dz^2 = dx_i dx_i \quad dS^2 = dX^2 + dY^2 + dZ^2 = dX_i dX_i$$



変位とひずみ

$$dX = \left(\frac{\partial X}{\partial x}\right) dx + \left(\frac{\partial X}{\partial y}\right) dy + \left(\frac{\partial X}{\partial z}\right) dz = \left(1 + \frac{\partial u}{\partial x}\right) dx + \left(\frac{\partial u}{\partial y}\right) dy + \left(\frac{\partial u}{\partial z}\right) dz$$

$$dY = \left(\frac{\partial Y}{\partial x}\right) dx + \left(\frac{\partial Y}{\partial y}\right) dy + \left(\frac{\partial Y}{\partial z}\right) dz = \left(\frac{\partial v}{\partial x}\right) dx + \left(1 + \frac{\partial v}{\partial y}\right) dy + \left(\frac{\partial v}{\partial z}\right) dz$$

$$dZ = \left(\frac{\partial Z}{\partial x}\right) dx + \left(\frac{\partial Z}{\partial y}\right) dy + \left(\frac{\partial Z}{\partial z}\right) dz = \left(\frac{\partial w}{\partial x}\right) dx + \left(\frac{\partial w}{\partial y}\right) dy + \left(1 + \frac{\partial w}{\partial z}\right) dz$$

$$\begin{aligned} dS^2 - ds^2 &= \left[\left(1 + \frac{\partial u}{\partial x}\right) dx + \left(\frac{\partial u}{\partial y}\right) dy + \left(\frac{\partial u}{\partial z}\right) dz \right]^2 - dx^2 \\ &+ \left[\left(\frac{\partial v}{\partial x}\right) dx + \left(1 + \frac{\partial v}{\partial y}\right) dy + \left(\frac{\partial v}{\partial z}\right) dz \right]^2 - dy^2 \\ &+ \left[\left(\frac{\partial w}{\partial x}\right) dx + \left(\frac{\partial w}{\partial y}\right) dy + \left(1 + \frac{\partial w}{\partial z}\right) dz \right]^2 - dz^2 \end{aligned}$$

変位とひずみ2

$$\begin{aligned} &= \left[2 \frac{\partial u}{\partial x} + \left(\frac{\partial u}{\partial x}\right)^2 + \left(\frac{\partial v}{\partial x}\right)^2 + \left(\frac{\partial w}{\partial x}\right)^2 \right] dx^2 \\ &+ \left[2 \frac{\partial v}{\partial y} + \left(\frac{\partial u}{\partial y}\right)^2 + \left(\frac{\partial v}{\partial y}\right)^2 + \left(\frac{\partial w}{\partial y}\right)^2 \right] dy^2 \\ &+ \left[2 \frac{\partial w}{\partial z} + \left(\frac{\partial u}{\partial z}\right)^2 + \left(\frac{\partial v}{\partial z}\right)^2 + \left(\frac{\partial w}{\partial z}\right)^2 \right] dz^2 \\ &+ \left[\left(\frac{\partial u}{\partial y}\right) + \left(\frac{\partial v}{\partial x}\right) + \left(\frac{\partial u}{\partial x}\right) \left(\frac{\partial u}{\partial y}\right) + \left(\frac{\partial v}{\partial x}\right) \left(\frac{\partial v}{\partial y}\right) + \left(\frac{\partial w}{\partial x}\right) \left(\frac{\partial w}{\partial y}\right) \right] dx dy \\ &+ \left[\left(\frac{\partial v}{\partial z}\right) + \left(\frac{\partial w}{\partial y}\right) + \left(\frac{\partial u}{\partial y}\right) \left(\frac{\partial u}{\partial z}\right) + \left(\frac{\partial v}{\partial y}\right) \left(\frac{\partial v}{\partial z}\right) + \left(\frac{\partial w}{\partial y}\right) \left(\frac{\partial w}{\partial z}\right) \right] dy dz \\ &+ \left[\left(\frac{\partial w}{\partial x}\right) + \left(\frac{\partial u}{\partial z}\right) + \left(\frac{\partial u}{\partial z}\right) \left(\frac{\partial u}{\partial x}\right) + \left(\frac{\partial v}{\partial z}\right) \left(\frac{\partial v}{\partial x}\right) + \left(\frac{\partial w}{\partial z}\right) \left(\frac{\partial w}{\partial x}\right) \right] dz dx \end{aligned}$$

ひずみ (Greenのひずみ)

$$\varepsilon_{xx} = \frac{\partial u}{\partial x} + \frac{1}{2} \left(\frac{\partial u}{\partial x}\right)^2 + \frac{1}{2} \left(\frac{\partial v}{\partial x}\right)^2 + \frac{1}{2} \left(\frac{\partial w}{\partial x}\right)^2$$

$$\varepsilon_{yy} = \frac{\partial v}{\partial y} + \frac{1}{2} \left(\frac{\partial u}{\partial y}\right)^2 + \frac{1}{2} \left(\frac{\partial v}{\partial y}\right)^2 + \frac{1}{2} \left(\frac{\partial w}{\partial y}\right)^2$$

$$\varepsilon_{zz} = \frac{\partial w}{\partial z} + \frac{1}{2} \left(\frac{\partial u}{\partial z}\right)^2 + \frac{1}{2} \left(\frac{\partial v}{\partial z}\right)^2 + \frac{1}{2} \left(\frac{\partial w}{\partial z}\right)^2$$

$$\varepsilon_{xy} = \frac{1}{2} \left[\left(\frac{\partial u}{\partial y}\right) + \left(\frac{\partial v}{\partial x}\right) + \left(\frac{\partial u}{\partial x}\right) \left(\frac{\partial u}{\partial y}\right) + \left(\frac{\partial v}{\partial x}\right) \left(\frac{\partial v}{\partial y}\right) + \left(\frac{\partial w}{\partial x}\right) \left(\frac{\partial w}{\partial y}\right) \right]$$

$$\varepsilon_{yz} = \frac{1}{2} \left[\left(\frac{\partial v}{\partial z}\right) + \left(\frac{\partial w}{\partial y}\right) + \left(\frac{\partial u}{\partial y}\right) \left(\frac{\partial u}{\partial z}\right) + \left(\frac{\partial v}{\partial y}\right) \left(\frac{\partial v}{\partial z}\right) + \left(\frac{\partial w}{\partial y}\right) \left(\frac{\partial w}{\partial z}\right) \right]$$

$$\varepsilon_{zx} = \frac{1}{2} \left[\left(\frac{\partial w}{\partial x}\right) + \left(\frac{\partial u}{\partial z}\right) + \left(\frac{\partial u}{\partial z}\right) \left(\frac{\partial u}{\partial x}\right) + \left(\frac{\partial v}{\partial z}\right) \left(\frac{\partial v}{\partial x}\right) + \left(\frac{\partial w}{\partial z}\right) \left(\frac{\partial w}{\partial x}\right) \right]$$

直ひずみ

剪断ひずみ

金属系の材料では1%以下

変位-ひずみの関係(2次元、線形)

- 変位 $\{u_x, u_y\}$
- ひずみ $\{\varepsilon_x, \varepsilon_y, \gamma_{xy}\}$
- 変位が十分小さいとき、線形の関係

$$\begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix} = \begin{Bmatrix} \partial u_x / \partial x \\ \partial u_y / \partial y \\ \partial u_x / \partial y + \partial u_y / \partial x \end{Bmatrix} = \begin{bmatrix} \partial / \partial x & 0 \\ 0 & \partial / \partial y \\ \partial / \partial y & \partial / \partial x \end{bmatrix} \begin{Bmatrix} u_x \\ u_y \end{Bmatrix}$$

ベクトル表示 $\varepsilon = \nabla^s u$

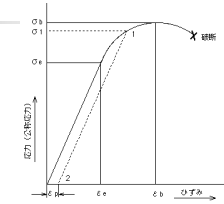
ひずみ-応力関係式

- Hookeの法則

$$\begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix} = \begin{bmatrix} D_{11} & D_{12} & D_{13} \\ D_{21} & D_{22} & D_{23} \\ D_{31} & D_{32} & D_{33} \end{bmatrix} \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix}$$

$$\begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix} = \frac{E}{1-\nu^2} \begin{bmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & (1-\nu)/2 \end{bmatrix} \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix}$$

$$\sigma = D\varepsilon$$

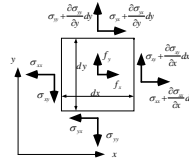


等方弾性
(平面応力問題)

E : ヤング率
 ν : ポアソン比

応力と外力の釣り合い

- 微小領域における釣り合い



$$\begin{Bmatrix} \frac{\partial \sigma_x}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + f_x \\ \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_y}{\partial y} + f_y \end{Bmatrix} = \begin{bmatrix} \partial / \partial x & 0 & \partial / \partial y \\ 0 & \partial / \partial y & \partial / \partial x \end{bmatrix} \begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix} + \begin{Bmatrix} f_x \\ f_y \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix}$$

$$\nabla^s T \sigma + f = 0$$

境界条件

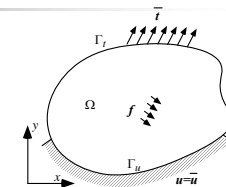
- 変位の境界条件

$$u = \bar{u}$$

- 荷重の境界条件

$$N_x \sigma = \bar{t}$$

$$N_x = \begin{bmatrix} n_x & 0 & n_y \\ 0 & n_y & n_x \end{bmatrix} \quad n = \begin{Bmatrix} n_x \\ n_y \end{Bmatrix}$$

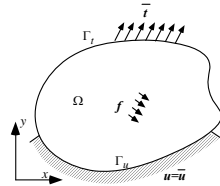


基礎方程式

$$\nabla^{*T}(D\nabla^*u) + f = 0$$

$$u = \bar{u} \text{ on } \Gamma_u$$

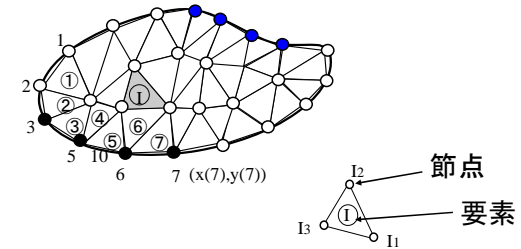
$$N_x D\nabla^*u = \bar{t} \text{ on } \Gamma_t$$



偏微分方程式の境界値問題
変位表記
状態変数 変位
自己随伴
解析解を求めるのは困難
→数値解

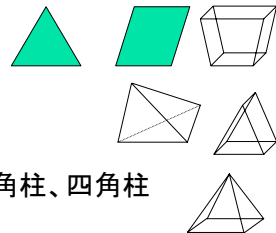
有限要素離散

- 領域を、単位領域(要素)に分割



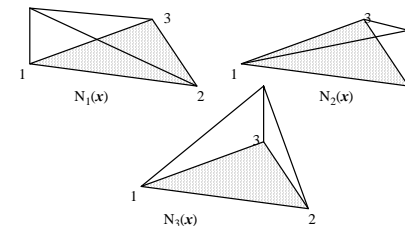
有限要素法のメッシュ

- 2次元
 - 三角形、四辺形
 - 1次要素、高次要素
- 3次元
 - 四面体、六面体、三角柱、四角柱
- プログラムによってはこれ以外のサポートもあるケースがある。

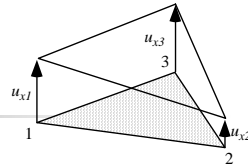


形状関数

- 要素内のある節点で1となり、他の節点で0となる関数
- 三角形要素では1次関数



変位の近似



■ 形状関数と節点での値

$$\tilde{u}_x(x, y) = u_{x1}N_1(x, y) + u_{x2}N_2(x, y) + u_{x3}N_3(x, y)$$

$$\tilde{u}_y(x, y) = u_{y1}N_1(x, y) + u_{y2}N_2(x, y) + u_{y3}N_3(x, y)$$

$$\tilde{\mathbf{u}}_i = \begin{Bmatrix} \tilde{u}_{xi} \\ \tilde{u}_{yi} \end{Bmatrix} = \begin{bmatrix} N_1 & 0 & N_2 & 0 & N_3 & 0 \\ 0 & N_1 & 0 & N_2 & 0 & N_3 \end{bmatrix} \begin{Bmatrix} u_{x1} \\ u_{y1} \\ u_{x2} \\ u_{y2} \\ u_{x3} \\ u_{y3} \end{Bmatrix}$$

$$\tilde{\mathbf{u}}_i = \mathbf{N}\mathbf{u}_i$$

ひずみ、応力の近似

- 変位を微分
- 三角形要素では要素内で一定値

$$\tilde{\boldsymbol{\varepsilon}}_i = \nabla^* \mathbf{u}_i = \nabla^* \mathbf{N}\mathbf{u}_i \equiv \mathbf{B}\mathbf{u}_i$$

$$\mathbf{B} = \begin{bmatrix} \partial N_1 / \partial x & 0 & \partial N_2 / \partial x & 0 & \partial N_3 / \partial x & 0 \\ 0 & \partial N_1 / \partial y & 0 & \partial N_2 / \partial y & 0 & \partial N_3 / \partial y \\ \partial N_1 / \partial y & \partial N_1 / \partial x & \partial N_2 / \partial y & \partial N_2 / \partial x & \partial N_3 / \partial y & \partial N_3 / \partial x \end{bmatrix}$$

$$\boldsymbol{\sigma}_i = \mathbf{D}\boldsymbol{\varepsilon}_i = \mathbf{D}\mathbf{B}\mathbf{u}_i$$

弱形式とエネルギー原理

■ 変分原理

$$\int_{\Omega} \delta \mathbf{u}^T (\nabla^{*T} (\mathbf{D}\nabla^* \mathbf{u}) + \mathbf{f}) d\Omega - \int_{\Gamma_t} \delta \mathbf{u}^T ((N_x \mathbf{D}\nabla^* \mathbf{u} - \bar{\mathbf{t}})) d\Gamma = 0$$

$$\mathbf{u} = \bar{\mathbf{u}} \quad \text{on } \Gamma_u$$



$$\int_{\Omega} ((\nabla^* \delta \mathbf{u})^T \mathbf{D}\nabla^* \mathbf{u} - \delta \mathbf{u}^T \mathbf{f}) d\Omega - \int_{\Gamma_t} \delta \mathbf{u}^T \bar{\mathbf{t}} d\Gamma = 0$$

$$\mathbf{u} = \bar{\mathbf{u}} \quad \text{on } \Gamma_u$$

仮想仕事の原理

仮想変位による内部エネルギーの増加分 = 外力の仕事

$$\int_{\Omega} (\nabla^* \delta \mathbf{u})^T \mathbf{D}\nabla^* \mathbf{u} d\Omega = \int_{\Omega} \delta \mathbf{u}^T \mathbf{f} d\Omega + \int_{\Gamma_t} \delta \mathbf{u}^T \bar{\mathbf{t}} d\Gamma$$

なぜ弱形式か

- 1次関数で近似するとひずみは要素間で不連続
- 2階微分は存在しない

$$\int_{\Omega} ((\nabla^* \delta \mathbf{u})^T \mathbf{D}\nabla^* \mathbf{u} - \delta \mathbf{u}^T \mathbf{f}) d\Omega - \int_{\Gamma_t} \delta \mathbf{u}^T \bar{\mathbf{t}} d\Gamma = 0$$

要素ごとの積分

$$\sum_i \int_{\Omega_i} ((\nabla^* \delta \mathbf{u})^T \mathbf{D}\nabla^* \mathbf{u} - \delta \mathbf{u}^T \mathbf{f}) d\Omega - \sum_i \int_{\Gamma_{t_i}} \delta \mathbf{u}^T \bar{\mathbf{t}} d\Gamma = 0$$

有限要素離散の弱形式

- 変位 $u = Bu_i$
- 仮想変位 $\delta u = B\delta u_i$

$$\sum_i \int_{\Omega_i} ((B\delta u_i)^T DBu_i) d\Omega = \sum_i \left[\int_{\Omega_i} (N\delta u_i)^T f d\Omega + \int_{\Gamma_n} (N\delta u_i)^T \bar{t} d\Gamma \right]$$

$$\sum_i \int_{\Omega_i} (\delta u_i^T B^T DBu_i) d\Omega = \sum_i \left[\int_{\Omega_i} \delta u_i^T N^T f d\Omega + \int_{\Gamma_n} \delta u_i^T N^T \bar{t} d\Gamma \right]$$

要素剛性マトリクスと 全体剛性マトリクス

$$\sum_i \int_{\Omega_i} (\delta u_i^T B^T DBu_i) d\Omega = \sum_i \left[\int_{\Omega_i} \delta u_i^T N^T f d\Omega + \int_{\Gamma_n} \delta u_i^T N^T \bar{t} d\Gamma \right]$$

$$\delta u^T Ku = \delta u^T f$$

全体剛性マトリクス $K = \sum_i \int_{\Omega_i} (B^T DB) d\Omega$ 要素剛性マトリクス

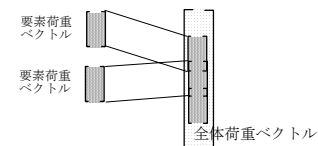
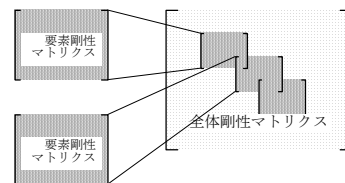
全体荷重ベクトル $f = \sum_i \left[\int_{\Omega_i} N^T f d\Omega + \int_{\Gamma_n} N^T \bar{t} d\Gamma \right]$ 要素荷重ベクトル

$$Ku = f$$

有限要素法の剛性方程式
→線形連立方程式ソルバー

剛性マトリクス、荷重ベクトルの 足し合わせ

- 全体自由度の中の相当する位置に足し込む



課題

- 板書で例題を行う
- 課題: 以下の問題を有限要素法を用いて解く。板書の続きを行い、要素剛性マトリクス、全体剛性マトリクスを求めよ。

